

# N-PENDULUM

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## Abstract

In this article we derive equations of motion for N-pendulum (in gravitational field and there is no other forces.), and study the numerical methods for solving these equations.

## I. Theory of N-pendulum

Consider  $N$  connected pendulums with lengths  $l_0, l_1, \dots, l_{N-1}$  and masses  $m_0, m_1, \dots, m_{N-1}$ . Position and velocity of  $n^{\text{th}}$  mass,  $\mathbf{r}_n, \mathbf{v}_n$ , relative to the suspending point is given by,

$$\mathbf{r}_n = \left( \sum_{i=0}^n l_i \sin q_i \right) \hat{\mathbf{x}} - \left( \sum_{i=0}^n l_i \cos q_i \right) \hat{\mathbf{y}}, \quad (1)$$

$$\mathbf{v}_n = \left( \sum_{i=0}^n l_i \dot{q}_i \cos q_i \right) \hat{\mathbf{x}} + \left( \sum_{i=0}^n l_i \dot{q}_i \sin q_i \right) \hat{\mathbf{y}}, \quad (2)$$

where  $q_i$  is the angle between the  $i^{\text{th}}$  pendulum and the vertical line. The kinetic and potential energy can be written as,

$$K = \sum_{n=0}^{N-1} \frac{1}{2} m_n \mathbf{v}_n \cdot \mathbf{v}_n = \sum_{n=0}^{N-1} \frac{1}{2} m_n \sum_{i=0}^n \sum_{j=0}^n l_i l_j \dot{q}_i \dot{q}_j \cos(q_i - q_j), \quad (3)$$

$$U = U_0 + \sum_{n=0}^{N-1} m g \hat{\mathbf{y}} \cdot \mathbf{r}_n = g \sum_{n=0}^{N-1} m_n \sum_{i=0}^n l_i (1 - \cos q_i), \quad (4)$$

and  $\mathcal{L} = K - U$ .

Now we can form equations of motion from Lagrange equations,

$$\frac{\partial \mathcal{L}}{\partial q_k} = -g l_k \sin q_k \sum_{n=k}^{N-1} m_n - \sum_{n=k}^{N-1} m_n \sum_{i=0}^n l_i l_k \dot{q}_i \sin(q_k - q_i), \quad (5)$$

$$\frac{\partial \mathcal{L}}{\partial \dot{q}_k} = \sum_{n=k}^{N-1} m_n \sum_{i=0}^n l_i l_k \dot{q}_i \cos(q_i - q_k), \quad (6)$$

$$\frac{d}{dt} \frac{\partial \mathcal{L}}{\partial \dot{q}_k} = \sum_{n=k}^{N-1} m_n \sum_{i=0}^n l_i l_k [\ddot{q}_i \cos(q_k - q_i) - \dot{q}_i (\dot{q}_k - \dot{q}_i) \sin(q_k - q_i)], \quad (7)$$

finally equations of motion are given by,

$$g l_k \sin q_k \sum_{n=k}^{N-1} m_n + \sum_{n=k}^{N-1} m_n \sum_{i=0}^n l_i l_k [\ddot{q}_i \cos(q_k - q_i) + \dot{q}_i^2 \sin(q_k - q_i)] = 0, \quad k = 0, 1, \dots, N-1, \quad (8)$$

we can write this equations as  $(\mathbf{q} = (q_1, \dots, q_N), \dot{\mathbf{q}} = (\dot{q}_1, \dots, \dot{q}_N))$ ,

$$\boxed{\mathbf{M}\ddot{\mathbf{q}} = \mathbf{F}(\mathbf{q}, \dot{\mathbf{q}})}, \quad (9)$$

$$M_{ki} := \left( \sum_{n=\max(k,i)}^{N-1} m_n \right) l_k l_i \cos(q_k - q_i), \quad (10)$$

$$F_k := -gl_k \sin q_k \sum_{n=k}^{N-1} m_n - \sum_{n=k}^{N-1} m_n \sum_{i=0}^n l_i l_k \dot{q}_i^2 \sin(q_k - q_i). \quad (11)$$

Also generalized momentum,  $p_k$ , can be written as,

$$p_k := \frac{\partial \mathcal{L}}{\partial \dot{q}_k} = \sum_{n=k}^{N-1} m_n \sum_{i=0}^n l_i l_k \dot{q}_i \cos(q_i - q_k) = (\mathbf{M}\dot{\mathbf{q}})_k. \quad (12)$$

We can drive equations of motion from Hamiltonian by using equations,

$$\begin{aligned} \dot{p}_k &= -\frac{\partial H}{\partial q_k}, \\ \dot{q}_k &= +\frac{\partial H}{\partial p_k}. \end{aligned}$$

For this purpose we must write Hamiltonian in terms of  $q_i, p_i$  instead of  $q_i, \dot{q}_i$ ,

$$\mathcal{H} = \frac{1}{2} \mathbf{M}\dot{\mathbf{q}}^T \dot{\mathbf{q}} + U = \frac{1}{2} \mathbf{p}^T \mathbf{M}^{-1} \mathbf{p} + U, \quad (13)$$

by this canonical Hamiltonian we can drive the equations of motion in terms of  $(\mathbf{q}, \mathbf{p})$ . Nevertheless we use equation (9) and also (12) for numerical computing. In fact we compute  $(\mathbf{q}, \dot{\mathbf{q}})$  from equation (9) and in each time-step compute  $\mathbf{p}$  from  $\mathbf{p} = \mathbf{M}\dot{\mathbf{q}}$ .

## II. Numerical Method

The equations (8) are non-linear and can not be solved analytically. Here we describe a method for solving these equations numerically.

Take  $\mathbf{s} := (\mathbf{q}, \dot{\mathbf{q}}) := (q_0, \dots, q_{N-1}, \dot{q}_0, \dots, \dot{q}_{N-1})$  as state. The equations of motion can be written as,

$$\mathbf{M}(\mathbf{q}) \ddot{\mathbf{q}} = \mathbf{F}(\mathbf{q}, \dot{\mathbf{q}}) \quad (14)$$

$$\Rightarrow \dot{\mathbf{s}} = \begin{bmatrix} \dot{\mathbf{q}} \\ \mathbf{M}^{-1} \mathbf{F}(\mathbf{q}, \dot{\mathbf{q}}) \end{bmatrix} = \mathbf{f}(\mathbf{s}) \quad (15)$$

which  $\mathbf{M}$  is a state-dependent matrix and  $\mathbf{F}$  is non-linear function given from (8).

Now we can use different algorithms for solving  $\dot{\mathbf{s}} = \mathbf{f}(\mathbf{s})$ . Euler-Richardson algorithm is good choice for this case, because the function  $\mathbf{f}$  depends on velocity.